



# Control of hot carrier thermalization in type-II quantum wells: a route to practical hot carrier solar cells



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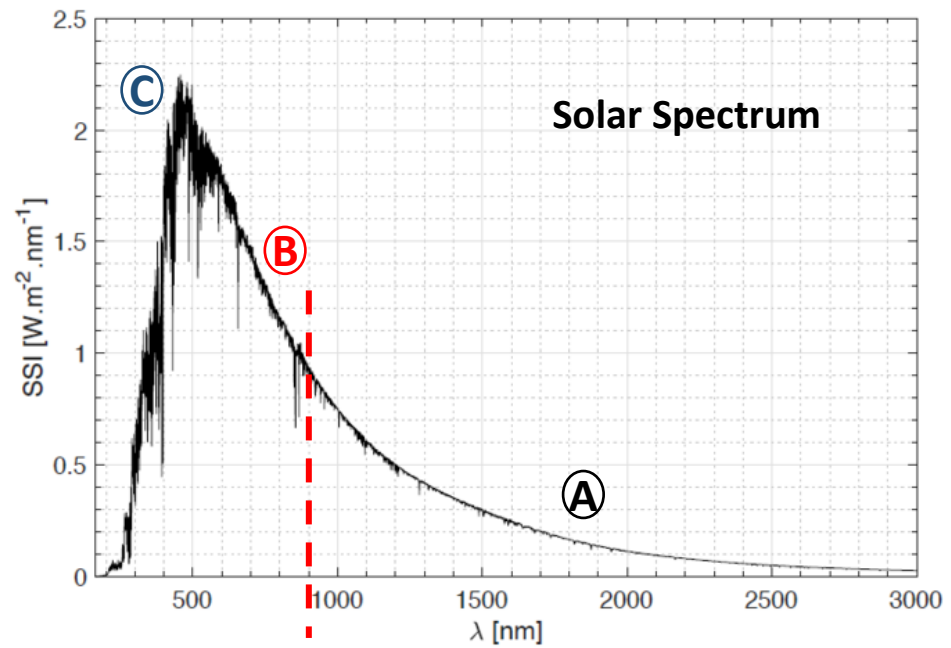
<sup>6</sup> The George Washington University, 21211 Street NW, Washington D. C. 20037



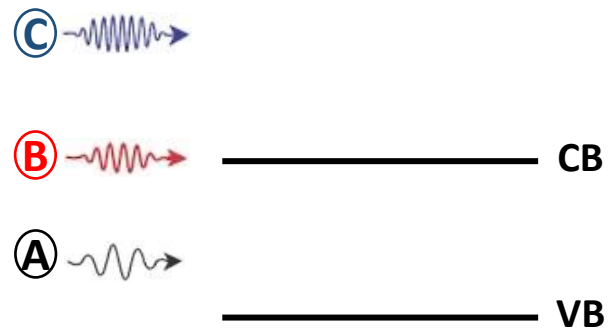
- Introduction: Hot Carrier Solar Cells
- InAs/AlAs<sub>x</sub>Sb<sub>1-x</sub> QWs and Optical Properties
- Hot Carrier Temperatures and Thermalization Coefficient
- Dynamics of Hot Carriers
- DFT Calculations and Raman Spectroscopy
- Electrical Characterization: Hot Carrier p-i-n diode



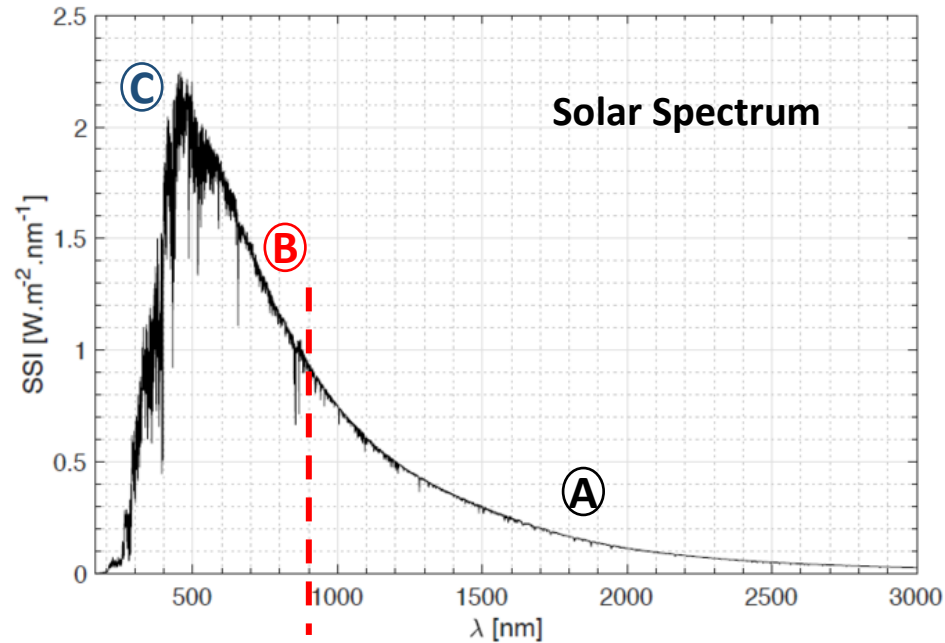
# Introduction: Hot Carrier Solar cell



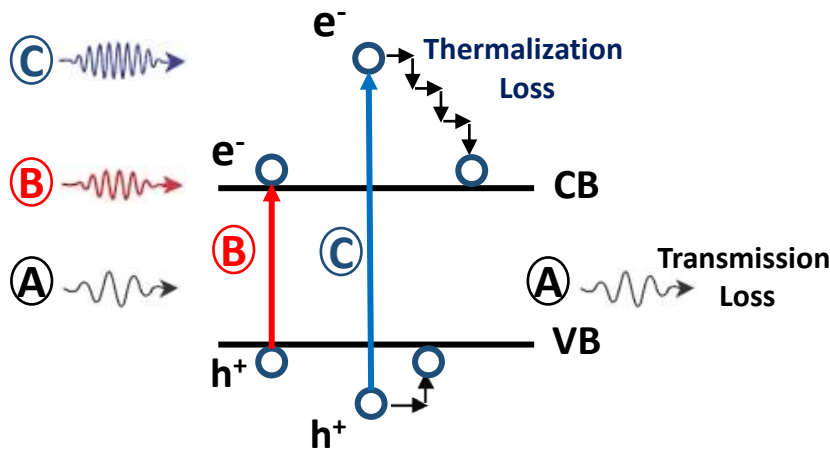
[https://www.esa.int/spaceinimages/Images/2017/12/Solar\\_spectrum](https://www.esa.int/spaceinimages/Images/2017/12/Solar_spectrum)



# Introduction: Hot Carrier Solar cell



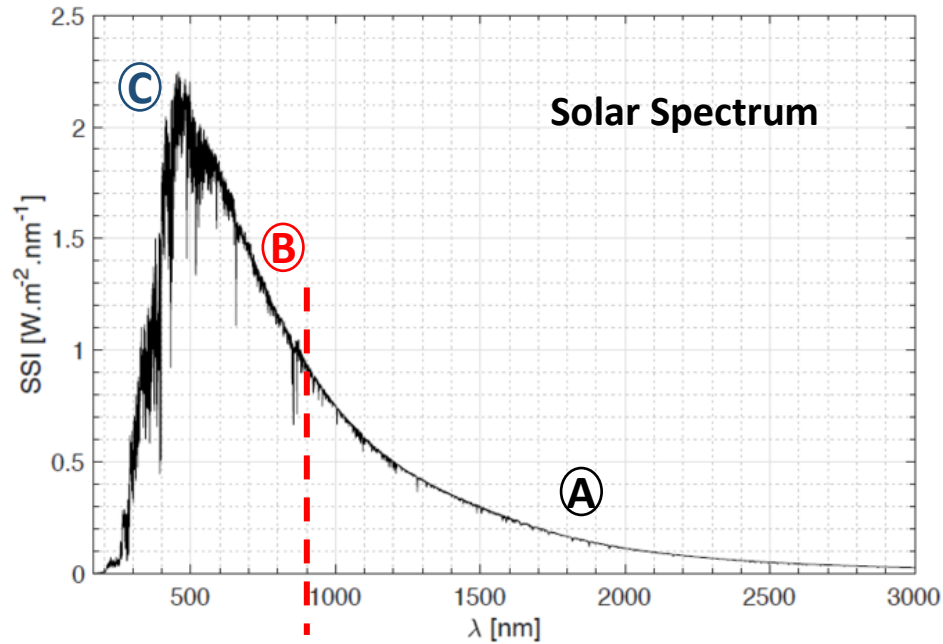
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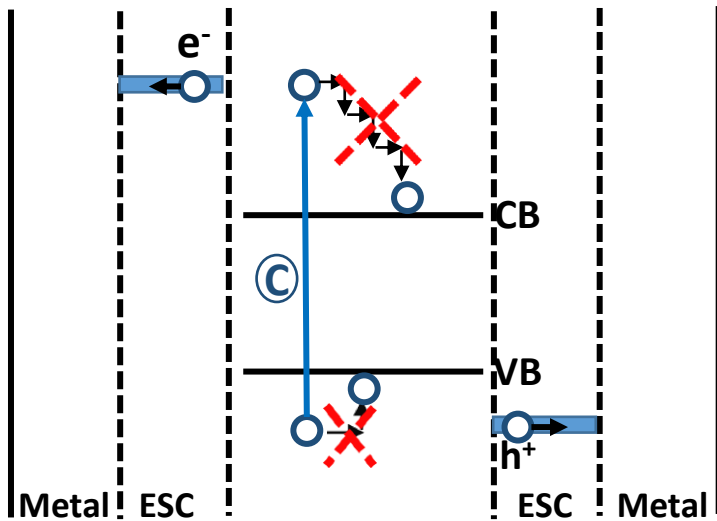
- **Maximum Theoretical Conversion Efficiency:**  
**33% No-Concentration (1 sun)**  
**(Shockley-Queisser Limit)**



# Introduction: Hot Carrier Solar cell



[https://www.esa.int/spaceinimages/Images/2017/12/Solar\\_spectrum](https://www.esa.int/spaceinimages/Images/2017/12/Solar_spectrum)



➤ **Maximum Theoretical Conversion Efficiency:**  
68% No-Concentration (1 sun)



# Determination of carrier temperature

## • Linear fit

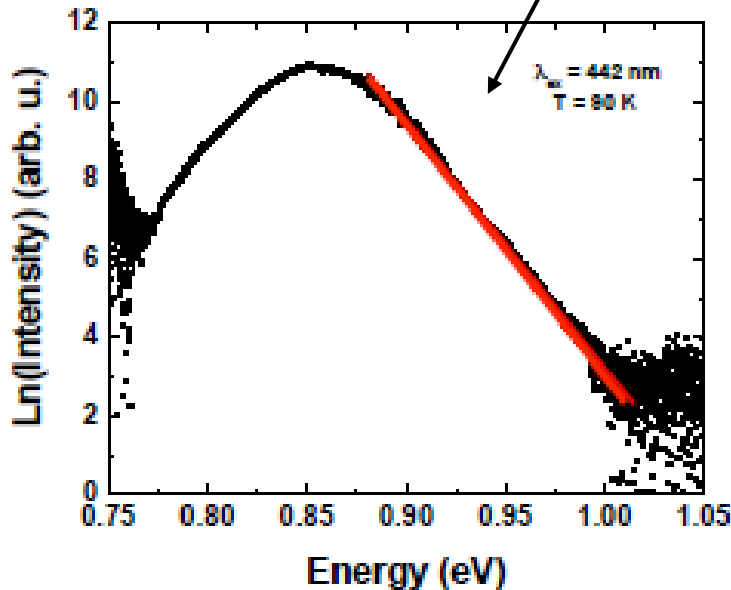


$$I_{PL}(E) = \frac{A(E) (E)^2}{4\pi^2 h^3 c^2} \left[ \exp\left(\frac{E - \Delta\mu}{k_B T}\right) - 1 \right]^{-1}$$

Generalized Planck's law

- Lasher & Stern, *Phys. Rev.* **133**, A553 (1964)
- De Vos & Pauwels, *Appl. Phys.* **25**, 119 (1981)
- P Würfel, *J. Phys. C: Solid State Phys.* **15** 3967 (1982)

$$\ln(I_{pl}(E)) \propto \left( -\frac{E}{k_B T_{eh}} \right)$$



### Related Research:

- A. Le Bris, L. Lombez, J. F. Guillemoles, *Energy & Environmental Science*, 5(3), 6225-6232 (2012).
- G. Conibeer, N. J. Ekins-Daukes, J. F. Guillemoles, D. König, M. Green, *Solar Energy Materials and Solar Cells*, 93(6-7), 713-719 (2009).
- L. C. Hirst, M. Sugiyama, N. J. Ekins-Daukes, *IEEE Journal of Photovoltaics*, 4(1), 244-252 (2014).



# Determination of carrier temperature

## • Linear fit

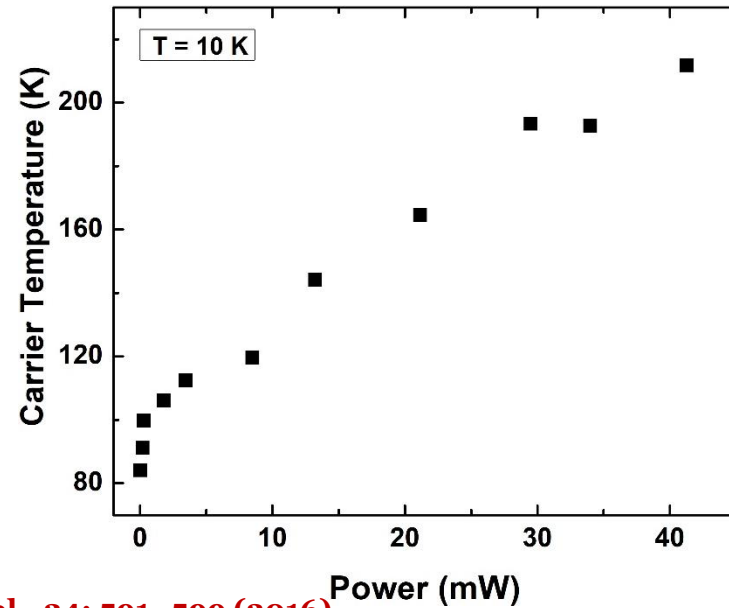
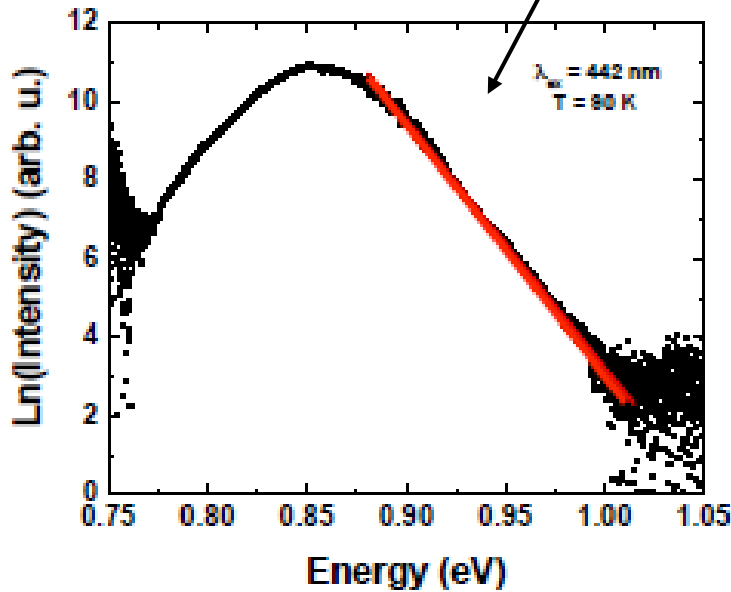


$$I_{PL}(E) = \frac{A(E) (E)^2}{4\pi^2 h^3 c^2} \left[ \exp\left(\frac{E - \Delta\mu}{k_B T}\right) - 1 \right]^{-1}$$

Generalized Planck's law

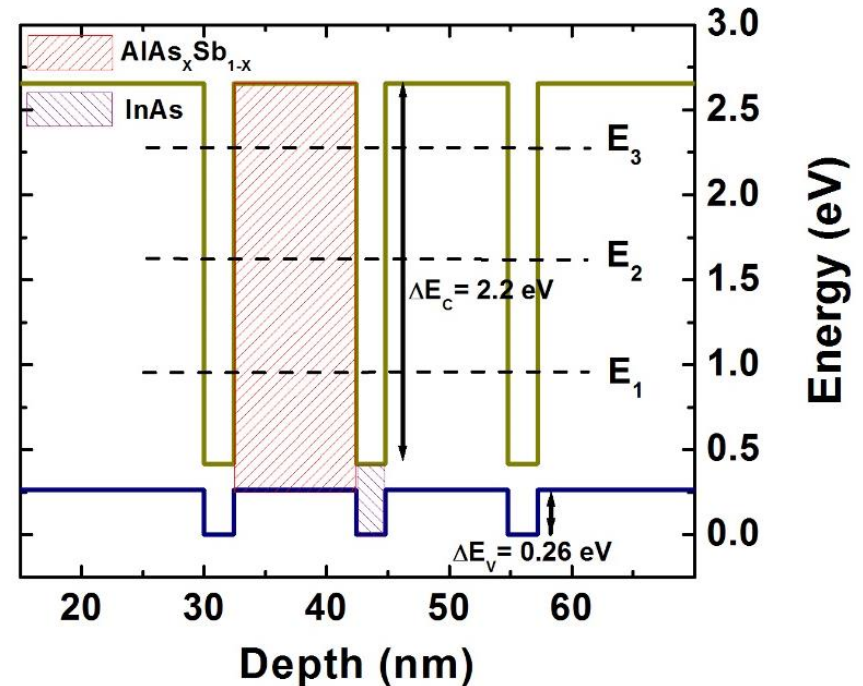
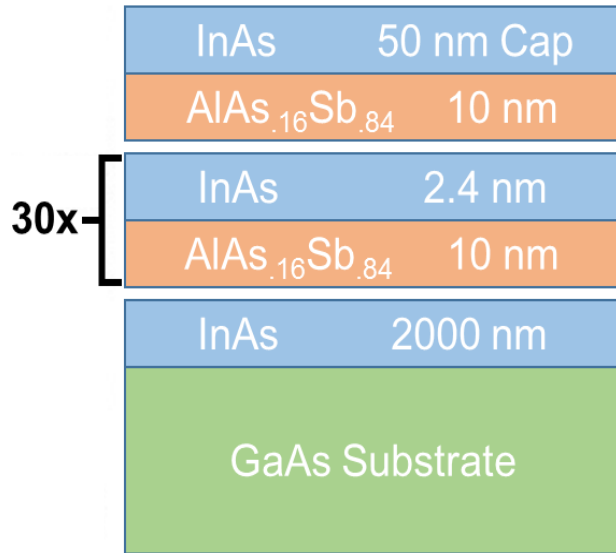
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$$\ln(I_{pl}(E)) \propto \left( -\frac{E}{k_B T_{eh}} \right)$$



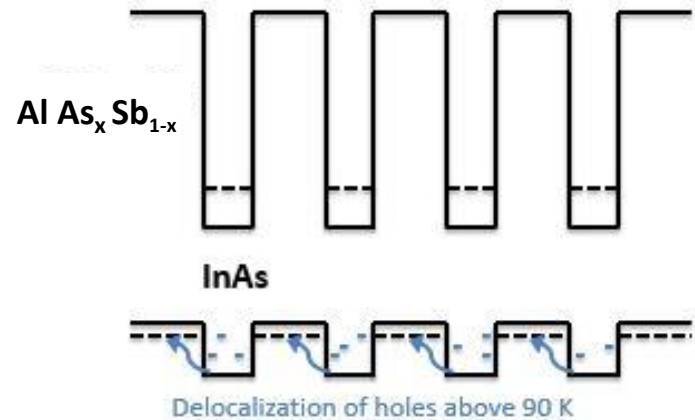
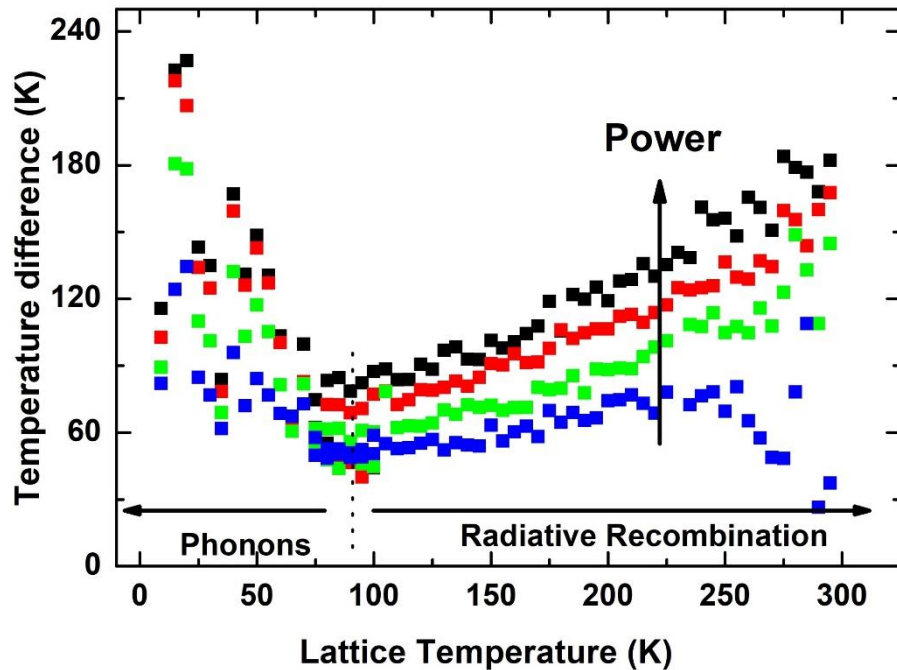
H. Esmailpour, I. R. Sellers, *et al.* *Prog. Photovolt: Res. Appl.*, **24**: 591–599 (2016).

# Type-II InAs/AlAs<sub>0.16</sub>Sb<sub>0.84</sub> quantum well structure



- “Deep” QW in the conduction band
- Quasi-type I structure
- Low quantum confinement in valence band
- Potential of creating resonant tunneling through a superlattice structure

**J. Tang, H. Esmailpour *et al.* Appl. Phys. Lett. 106, 061902 (2015)**



- ❖ Below 90 K, phonon mediated thermalization of hot carriers due to localized states at the interface is dominant.
- ❖ Above 90 K, radiative recombination of spatially separated photogenerated charges in type-II band alignment is the dominant mechanism which creates robust hot carrier effect at elevated temperatures.

**J. Tang, H. Esmailpour *et al.* Appl. Phys. Lett. 106, 061902 (2015)**





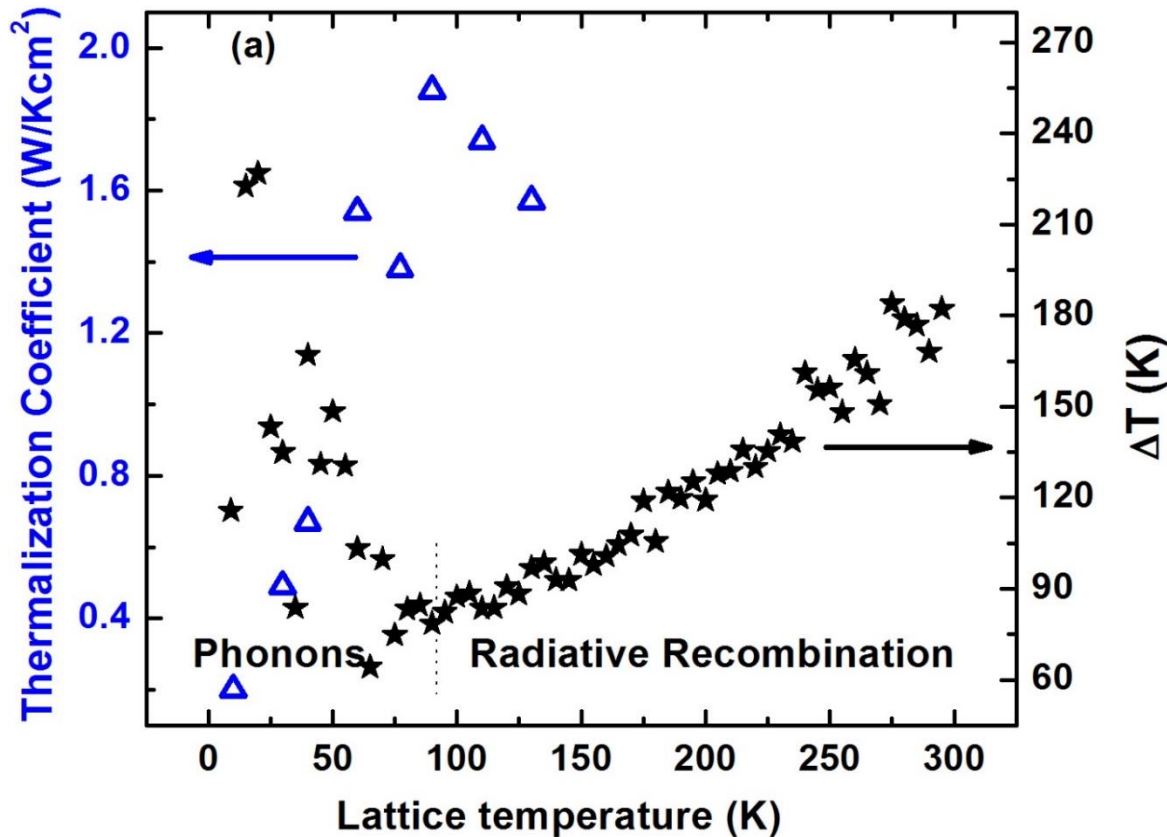
# Thermalization coefficient in InAs/AlAsSb QWs



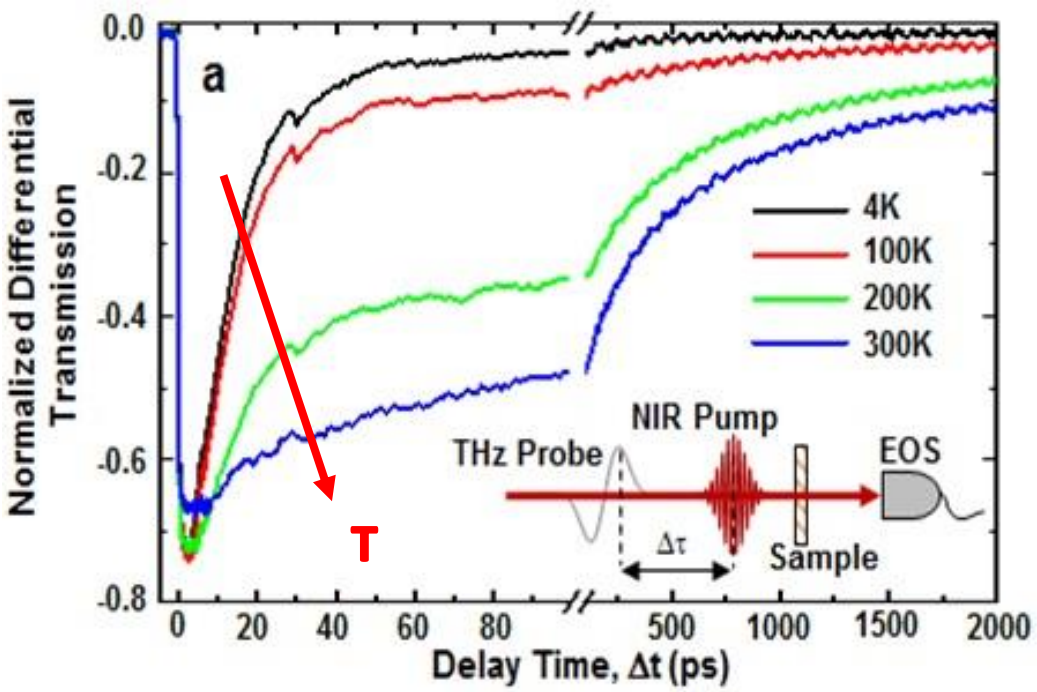
$$P_{abs} = P_{th} = Q(T_{eh} - T) \exp\left(-\frac{h\nu_{LO}}{k_B T_{eh}}\right)$$

Le Bris, A., *et al.* Energy & Environmental Science 5.3 (2012): 6225-6232.

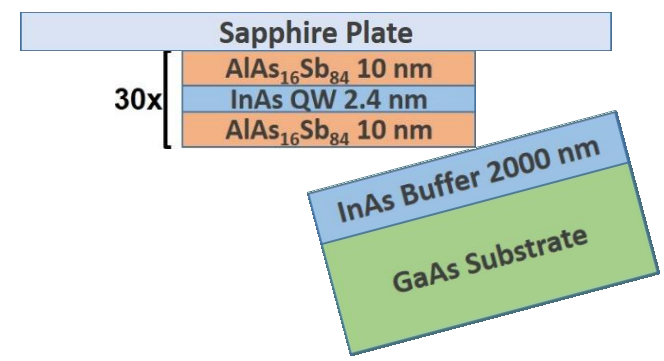
Q: Thermalization coefficient



H. Esmailpour, I. R. Sellers, *et al.* Prog. Photovolt: Res. Appl., 24: 591–599 (2016).



*Processed Sample for Ultra-fast Transient Absorption Spectroscopy*



$$\Delta E(t)/E = \sum_{i=1}^3 A_i \exp[-(t - t_0)/\tau_i]$$

*E*: THz Electric Fields  
*A<sub>i</sub>*: Normalized Amplitude  
*τ<sub>i</sub>*: Decay Time

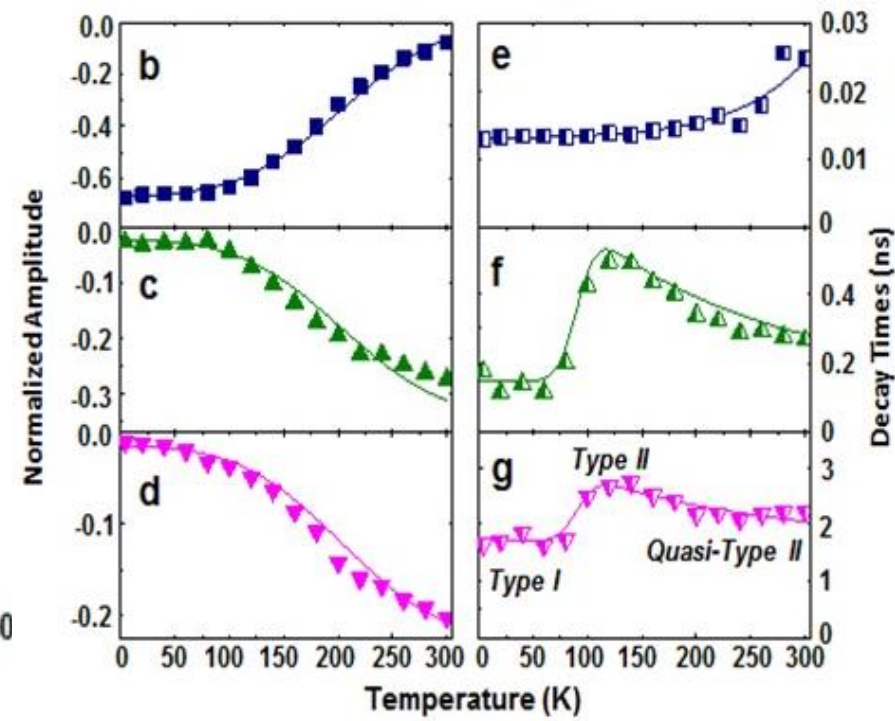
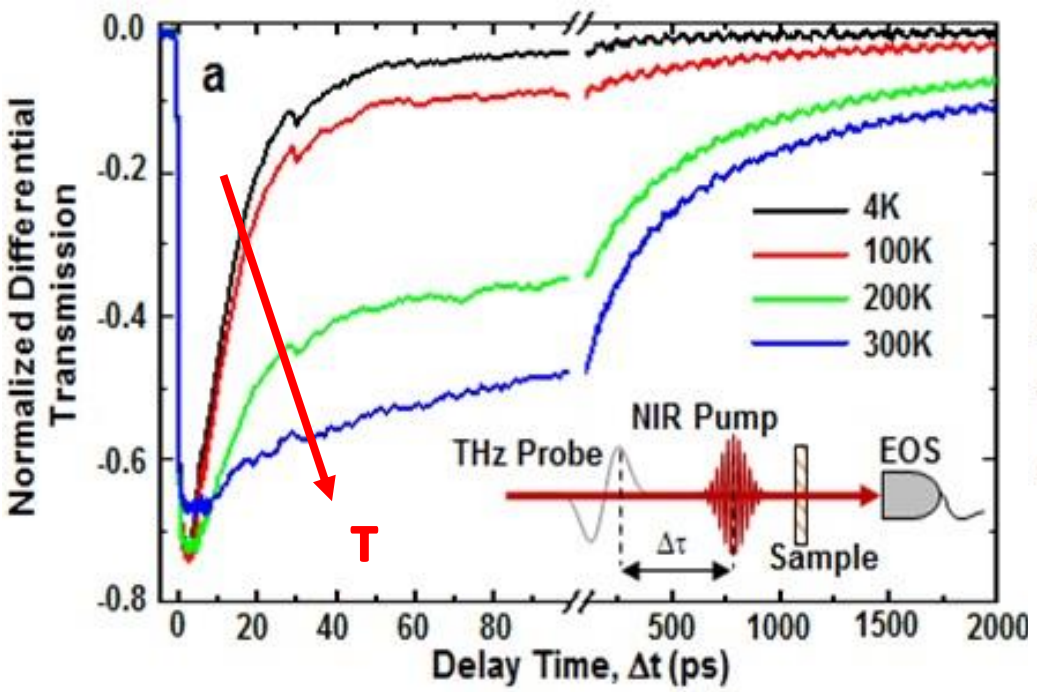


**H. P. Piyathilaka, A. D. Bristow**

**H. Esmailpour, I. R. Sellers, et al. Scientific Reports (under review).**



# Time Domain THz Spectroscopy: AC Photoconductivity



$$\Delta E(t)/E = \sum_{i=1}^3 A_i \exp[-(t - t_0)/\tau_i]$$

$E$ : THz Electric Fields  
 $A_i$ : Normalized Amplitude  
 $\tau_i$ : Decay Time



H. P. Piyathilaka, A. D. Bristow

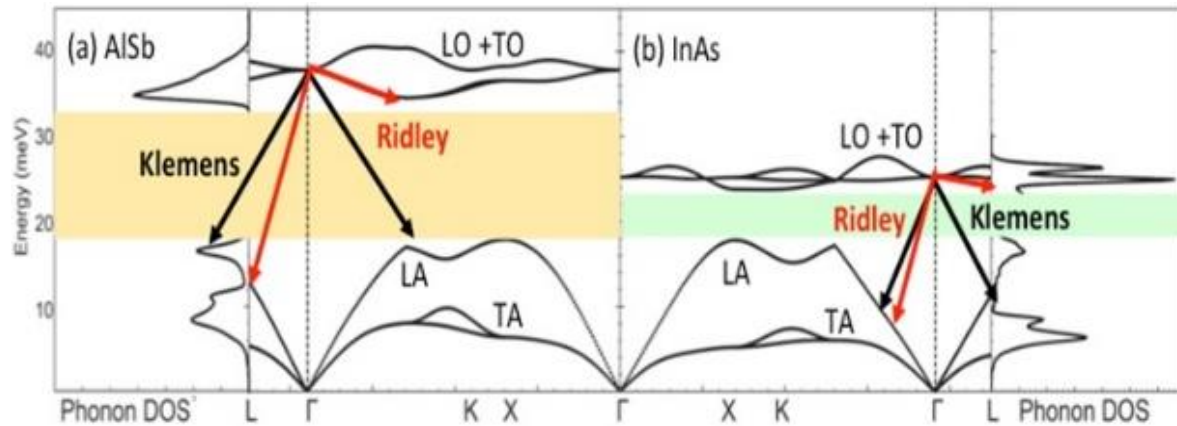
H. Esmailpour, I. R. Sellers, *et al.* Scientific Reports (under review).

V. R. Whiteside, H. Esmailpour, I. R. Sellers, *et al.* Semiconductor Science and Technology (under review).

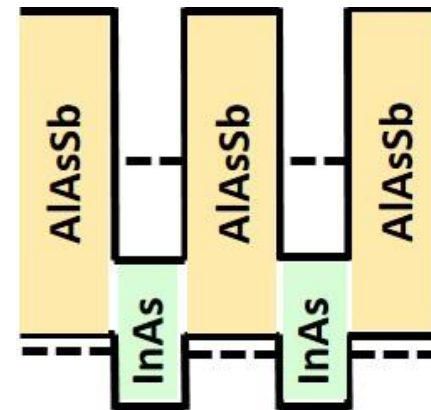


# Phononic Properties of InAs/AlAsSb QW Structure

- DFT Calculations



- ❖ Large optical phonon band gap for AlSb sample due to a large difference between cation and anion mass.
- ❖ Large difference between optic and acoustic phonon energy for AlSb ( $\hbar\omega_{LO}/\hbar\omega_{LA} \sim 1.9$ ) (see tan shaded region) and small difference between optic and acoustic phonon energy ( $\hbar\omega_{LO}/\hbar\omega_{LA} \sim 1.1$ ) for InAs (see green shaded region)



H. Esmailpour, I. R. Sellers, *et al.* Scientific Reports (under review).

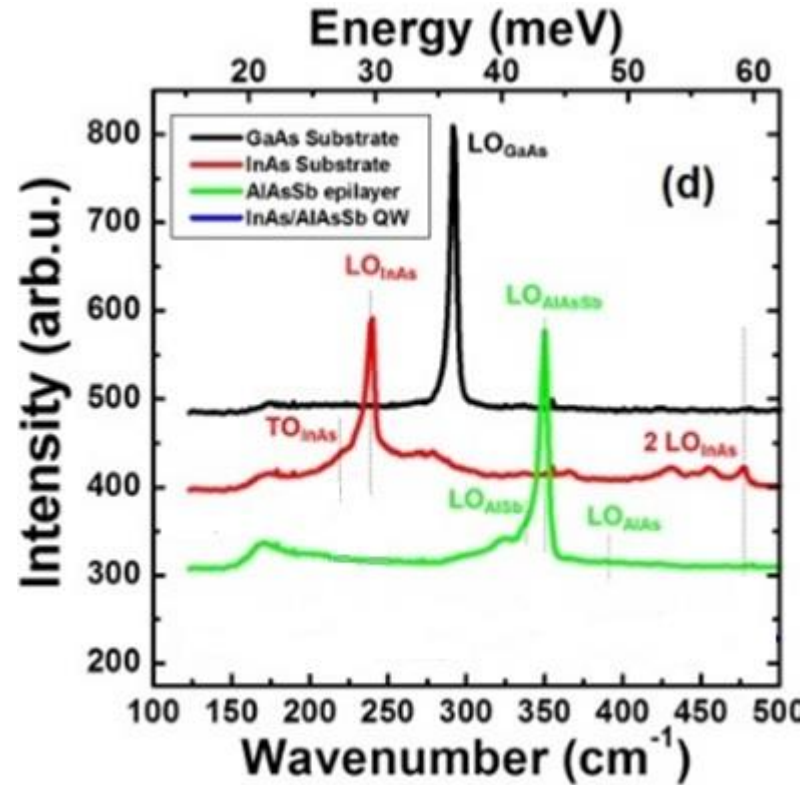
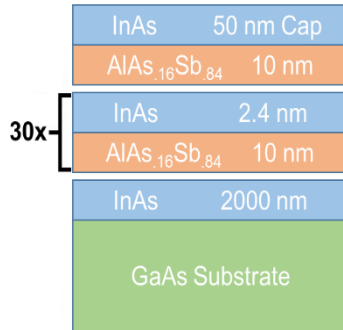


B. Wang



# Phononic Properties of InAs/AlAsSb QW Structure

- Raman Spectroscopy

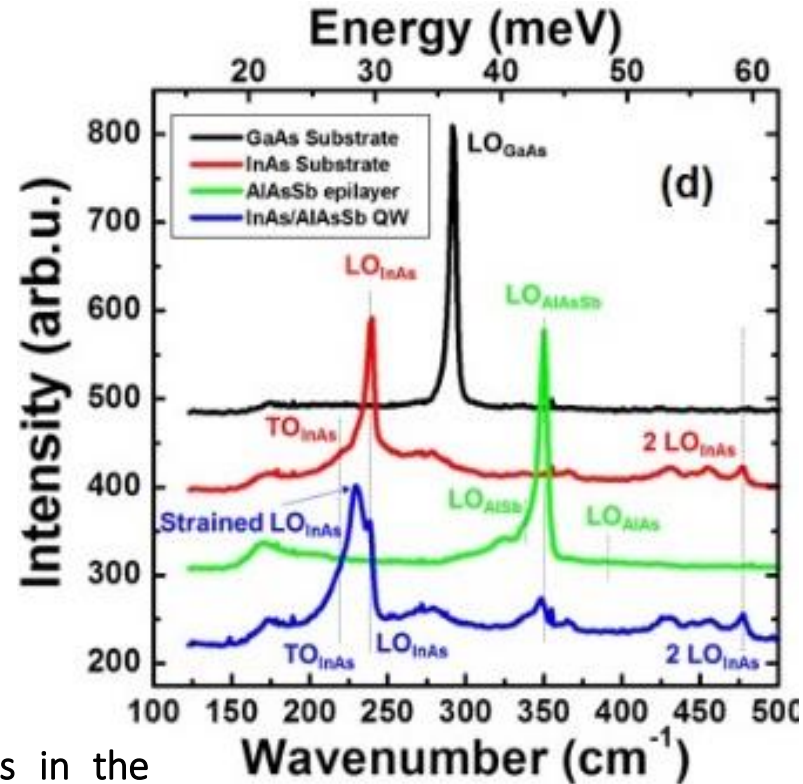
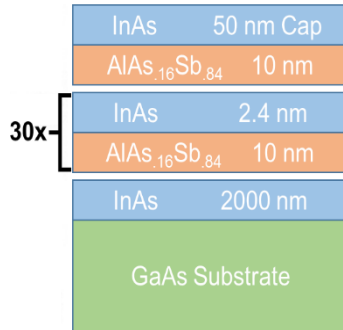


**E. Adcock-Smith,**  
**K. P. Roberts**

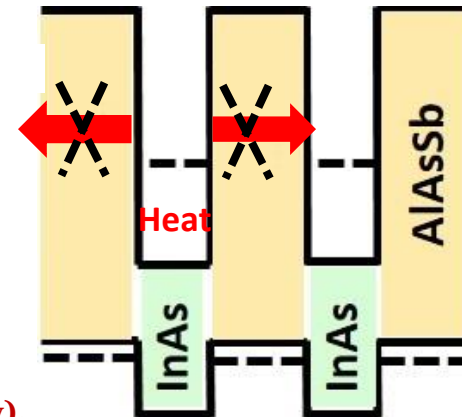
**H. Esmailpour, I. R. Sellers, et al. Scientific Reports (under review).**

Photovoltaics Materials & Device Group, University of Oklahoma: <http://www.nhn.ou.edu/~sellers/group/index.html>

- Raman Spectroscopy



- ❖ Limited contribution of phonon process in the barriers (AlAsSb) in carrier thermalization.
- ❖ Strong confinement of hot electron-phonon process within the QWs.
- ❖ Up-conversion of LA phonons to LO phonons and re-heating and stabilizing hot electron distribution within the InAs QW.



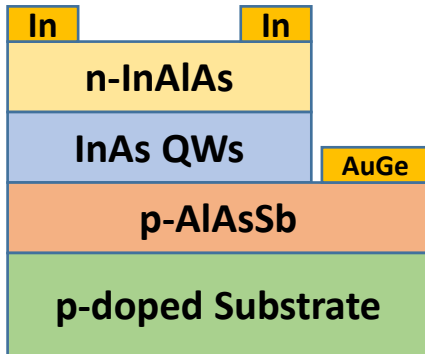
H. Esmailpour, I. R. Sellers, *et al.* Scientific Reports (under review).



# Electrical characterization: Hot carrier p-i-n diode

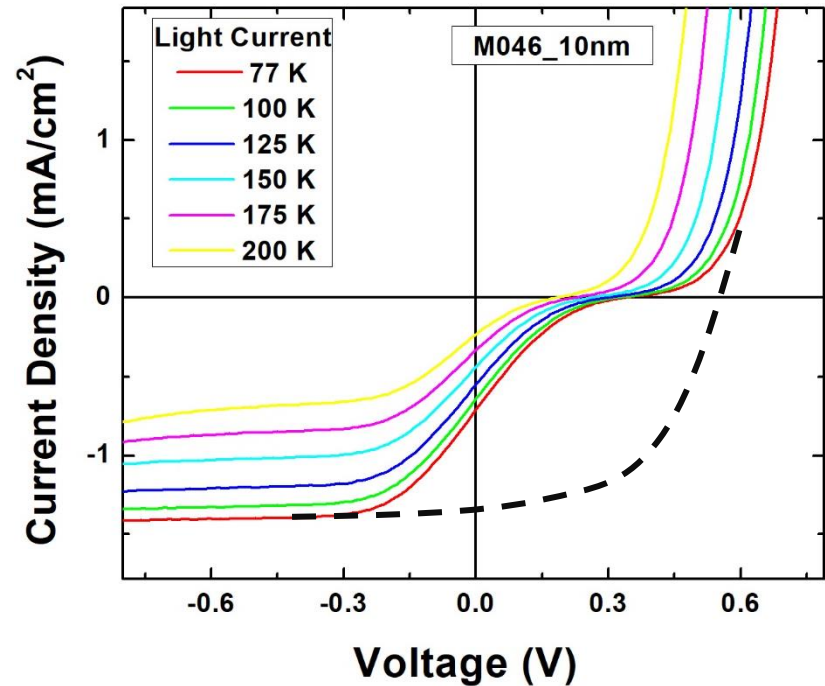


- I-V measurements as a function of barrier thickness



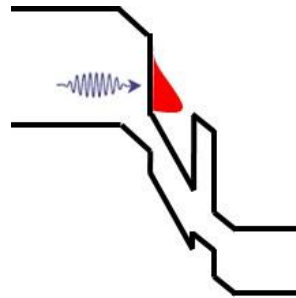
## Related Research:

- L. C. Hirst, R. J. Walters, M. F. Führer, N. J. Ekins-Daukes, *Applied Physics Letters*, 104(23), 231115, (2014).

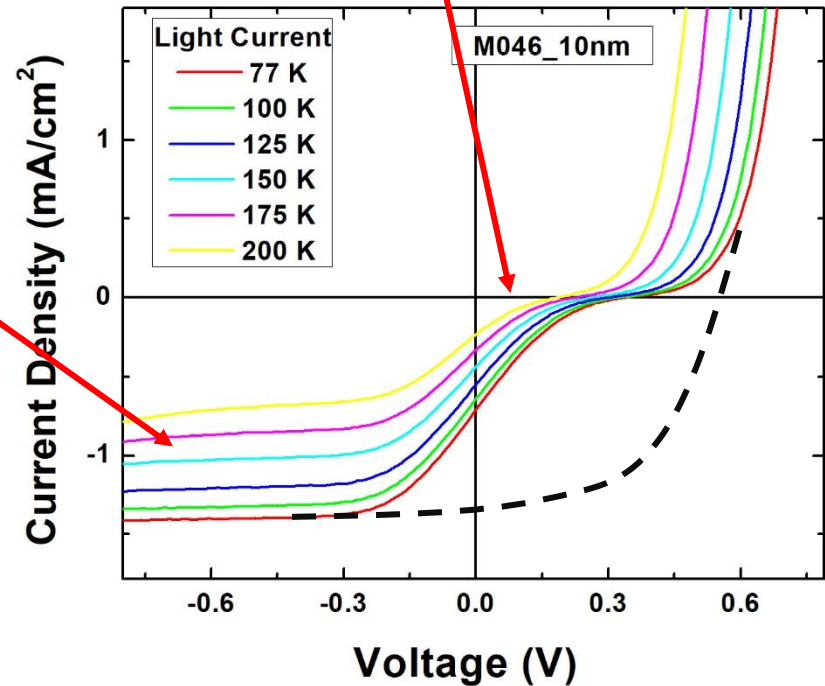
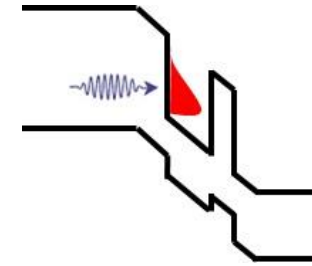


- I-V measurements as a function of barrier thickness

❖ At reverse bias where there is a large band bending, generated electrons can tunnel through the barrier material into the n-doped region.



❖ At forward bias, confinement of hot carriers within the QW was observed due to small band bending of the QW structure.

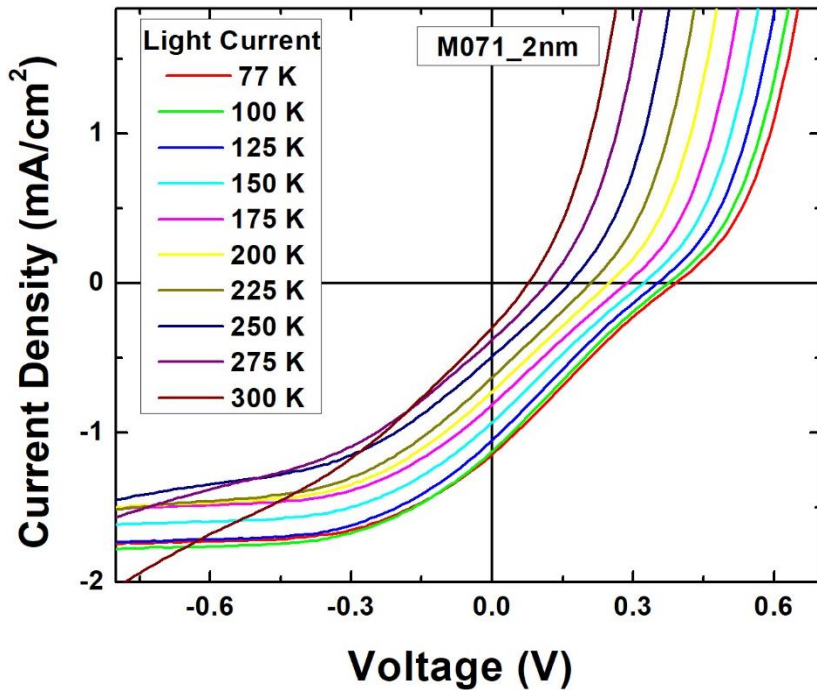




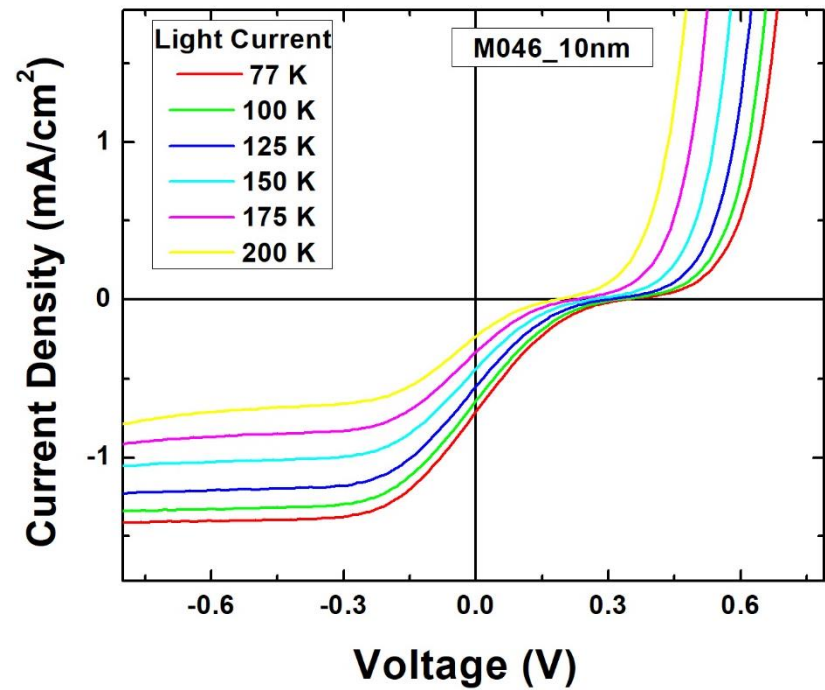
- I-V measurements as a function of barrier thickness

❖ Smaller carrier confinement within the QW is observed for the p-i-n device with a thinner barrier material.

### 2 nm barrier p-i-n diode



### 10 nm barrier p-i-n diode



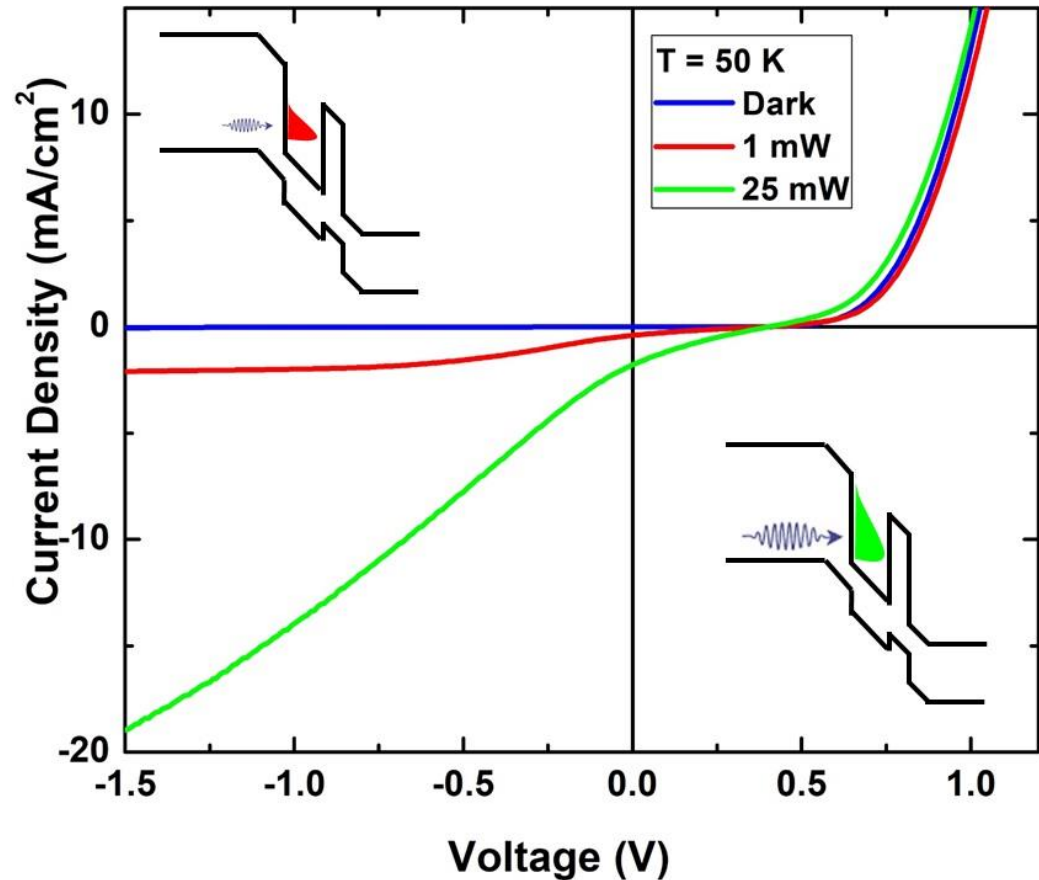


# Electrical characterization: Hot carrier p-i-n diode



- Power dependence

- ❖ By increasing the barrier thickness, less amount of current can escape from the confinement of the QW.
- ❖ As power increases, more hot electrons can escape from the QW.





# Summary

- InAs/Al<sub>x</sub>As<sub>1-x</sub>Sb quantum wells offer potential as active absorber in hot carrier solar cells.
- Thermalization factor analysis appears unsuitable for the proposed mechanism at higher temperatures.
- Evidence of state filling in the valence band has been observed.
- Carrier temperatures and chemical potentials for generated electrons and holes have been found using the non-equilibrium generalized Planck's law.
- Tera-hertz time domain spectroscopy (TDS) has been done using a processed sample made of only the InAs MQWs.
- Hot carrier p-i-n diodes have shown the existence of hot carriers confined within the InAs MQWs.

## Acknowledgments



ECCS # 1610062

